

LEY DE BIOT - SAVART

HOJA N°

FECHA

VECTORIAL :

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I d\vec{l} \times \hat{r}}{r^2}$$

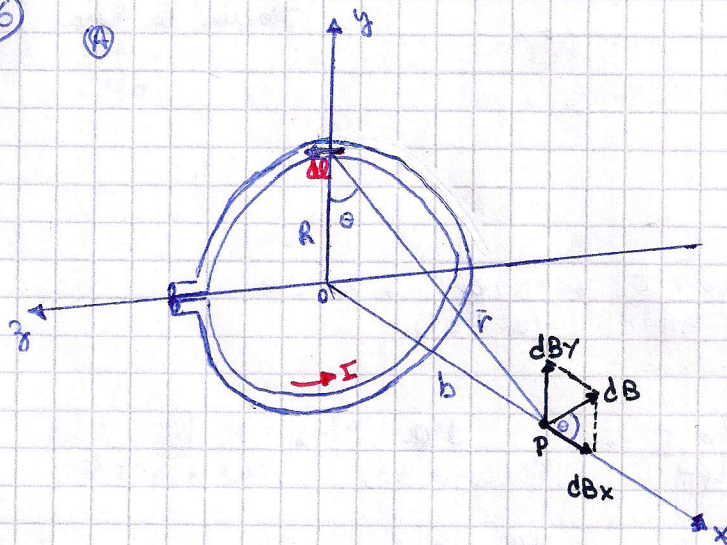
$$\mu_0 = 4\pi \cdot 10^{-7} \frac{T \cdot m}{A}$$

MÓDULO :

$$dB = \frac{\mu_0}{4\pi} \frac{I dl \cdot \sin \theta}{r^2}$$

• θ es el ángulo entre dl y r

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(B)

dl y $\vec{r}$ son perpendiculares $\Rightarrow \sin 90^\circ = 1$

$$dB = \frac{\mu_0}{4\pi} \cdot \frac{I \cdot dl}{r^2}$$

$$r^2 = b^2 + R^2$$

$$r = (b^2 + R^2)^{1/2}$$

$$dB = \frac{\mu_0}{4\pi} \cdot \frac{I dl}{(b^2 + R^2)}$$

Por simetría, $dBy = 0 \Rightarrow$ solo tenemos $dBx = dB \cdot \cos \theta$

$$\cos \theta = \frac{R}{(b^2 + R^2)^{1/2}} \Rightarrow dBx = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{dl}{(b^2 + R^2)} \cdot \frac{R}{(b^2 + R^2)^{1/2}}$$

$$dBx = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{dl}{(b^2 + R^2)^{3/2}} \cdot R$$

$$dl = R \cdot d\theta \Rightarrow dBx = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{R^2 d\theta}{(b^2 + R^2)^{3/2}}$$

$$B = Bx = \int dBx = \int \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{R^2 d\theta}{(b^2 + R^2)^{3/2}} = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{R^2}{(b^2 + R^2)^{3/2}} \int_0^{2\pi} d\theta$$

$$= \frac{\mu_0 I}{4\pi} \cdot \frac{R^2 2\pi}{(b^2 + R^2)^{3/2}} \Rightarrow$$

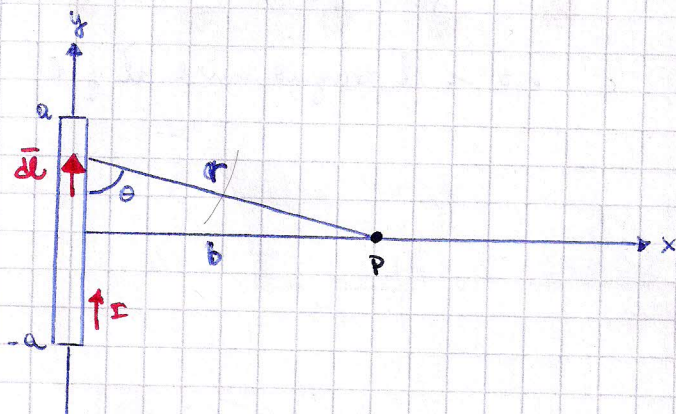
$$B = \frac{\mu_0 \cdot I}{2} \cdot \frac{R^2}{(b^2 + R^2)^{3/2}}$$

NOTA

© En el centro de la espira $\Rightarrow b = 0$

$$B = \frac{\mu_0 \cdot I}{2} \cdot \frac{R^2}{(0^2 + R^2)^{3/2}} \Rightarrow B = \frac{\mu_0 \cdot I}{2R}$$

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$dB = \odot$
Hacia la hoja

$$dl = da$$

$$dB = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{da}{(a^2 + b^2)} \cdot \sin \theta$$

$$dB = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{da}{(a^2 + b^2)} \cdot \frac{b}{(a^2 + b^2)^{1/2}}$$

$$dB = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{da \cdot b}{(a^2 + b^2)^{3/2}}$$

$$B = \int dB = \frac{\mu_0 \cdot I}{4\pi} \cdot \int_{-a}^a \frac{b}{(a^2 + b^2)^{3/2}} da \Rightarrow \text{POR TABLA}$$

$$B = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{2a}{b \cdot \sqrt{b^2 + a^2}}$$

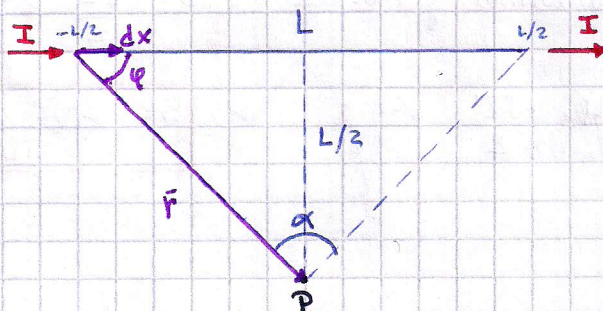
Como el conductor es infinito:

$$B = \lim_{a \rightarrow \infty} \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{2a}{b \cdot \sqrt{b^2 + a^2}}$$

$\xrightarrow{a}$

$$B = \frac{\mu_0 \cdot I}{2\pi b}$$

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$$\alpha = 90^\circ$$

$$r^2 = x^2 + \frac{L^2}{4}$$

$$\sin \phi = \frac{L/2}{\left(x^2 + \frac{L^2}{4}\right)^{1/2}}$$

$$dB = \frac{\mu_0 \cdot I}{4\pi} \cdot \frac{dx}{x^2 + \frac{L^2}{4}} \cdot \frac{L/2}{\left(x^2 + \frac{L^2}{4}\right)^{1/2}}$$

$$B = \int_{-L/2}^{L/2} dB = \frac{\mu_0 \cdot I}{4\pi} \cdot L/2 \int_{-L/2}^{L/2} \frac{dx}{\left(x^2 + \frac{L^2}{4}\right)^{3/2}} = \frac{\mu_0 I}{4\pi} \cdot L/2 \left[\frac{x}{\frac{L^2}{4} \left(x^2 + \frac{L^2}{4}\right)^{1/2}} \right]_{-L/2}^{L/2}$$

$$= \frac{\mu_0 I}{4\pi} \cdot \frac{L}{2} \cdot \frac{L}{L^2} \cdot \left[\frac{L/2}{\left(\frac{L^2}{4} + \frac{L^2}{4}\right)^{1/2}} - \frac{-L/2}{\left(\frac{L^2}{4} + \frac{L^2}{4}\right)^{1/2}} \right]$$

$$= \frac{\mu_0 \cdot I}{2\pi L} \cdot \left[\frac{L}{\left(\frac{L^2}{2}\right)^{1/2}} \right] = \frac{\mu_0 I}{2\pi L} \cdot \frac{L}{\frac{L}{\sqrt{2}}} = \frac{\mu_0 I}{2\pi L} \cdot \sqrt{2}$$

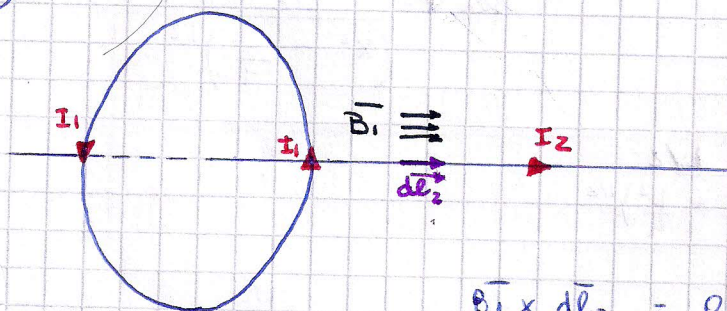
$$B = \frac{\mu_0 \cdot I}{\pi L} \cdot \frac{\sqrt{2}}{2}$$

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Para un solo alambre calculamos y nos dio: $B = \frac{\sqrt{z}}{z} \frac{\mu_0 \cdot I}{\pi L}$

Entonces para 4 sería: $4 \cdot \frac{\sqrt{z}}{z} \frac{\mu_0 I}{\pi L} \Rightarrow B = \frac{4\sqrt{z}}{z} \frac{\mu_0 \cdot I}{\pi L}$

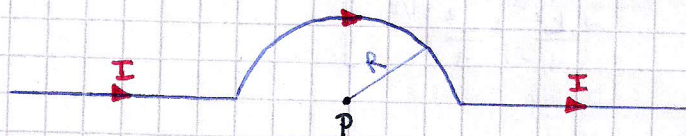
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$\vec{B}_1 \times d\vec{l}_2 = 0$; son paralelos

$$F_{22} = 0$$

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Regla de la mano derecha:

campo en P entrante $\odot$.

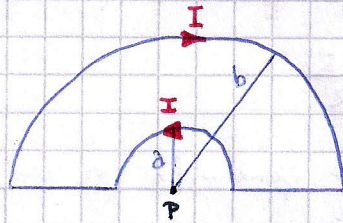
En el 162, el campo en el centro de la esfera dio:

$$B = \frac{\mu_0 \cdot I}{2R}$$

Entonces para la mitad sería $\frac{B}{2} \Rightarrow B = \frac{\mu_0 \cdot I}{4R}$

Los tramos rectos no contribuyen, ya que $d\vec{l}$ y $\vec{r}$ son paralelos.

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Regla de la mano derecha.

El campo según b en el centro será entrante ($-\hat{k}$) y según a será saliente ($+\hat{k}$).

$$\vec{B}_{pb} = \frac{\mu_0 \cdot I}{4b} (-\hat{k})$$

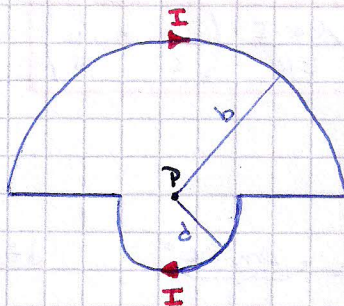
$$\vec{B}_{pa} = \frac{\mu_0 \cdot I}{4a} (\hat{k})$$

$$\Rightarrow \vec{B}_p = \vec{B}_{pb} + \vec{B}_{pa} = (\hat{k}) \cdot \left(\frac{\mu_0 \cdot I}{4a} - \frac{\mu_0 \cdot I}{4b} \right)$$

$$\vec{B}_p = \frac{\mu_0 \cdot I}{4} \hat{k} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$B = \frac{\mu_0 \cdot I}{4} \frac{b-a}{a \cdot b} \quad \text{SALIENTE}$$

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Regla de la mano derecha.

El campo según a y b será entrante ($-\hat{k}$).

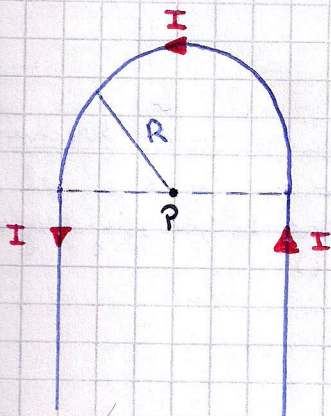
$$\vec{B}_{pb} = \frac{\mu_0 \cdot I}{4b} (-\hat{k})$$

$$\Rightarrow \vec{B}_p = \vec{B}_{pa} + \vec{B}_{pb} = (-\hat{k}) \frac{\mu_0 \cdot I}{4} \left(\frac{1}{a} + \frac{1}{b} \right)$$

$$\vec{B}_{pa} = \frac{\mu_0 \cdot I}{4a} (-\hat{k})$$

$$B = \frac{\mu_0 \cdot I}{4} \cdot \frac{a+b}{a \cdot b} \quad \text{ENTRANTE}$$

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El campo generado por la semi:

$$B = \frac{\mu_0 I}{4R}$$

El campo generado por un conductor infinito:

$$B = \frac{\mu_0 \cdot I}{2\pi R}$$

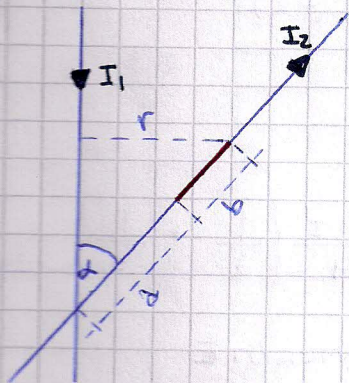
Como en el conductor contribuye infinitamente hacia abajo, podemos considerar solo la contribución de la mitad:

$$B = \frac{\mu_0 \cdot I}{4\pi R}$$

Entonces:

$$B_p = 2 \cdot \frac{\mu_0 \cdot I}{4\pi R} + \frac{\mu_0 \cdot I}{4R} \Rightarrow B = \frac{\mu_0 \cdot I}{2R} \left(\frac{1}{2} + \frac{1}{\pi} \right)$$

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$$B_{z1} = \frac{\mu_0 \cdot I_1}{2\pi r} \cdot \underbrace{\sin 90^\circ}_{=1}$$

$$dF_{z1} = I_2 \cdot dz \cdot B_{z1}$$

$$B_{z1} = \frac{\mu_0 \cdot I_1}{2\pi r}$$

$$dF_{z1} = I_2 \cdot dz \cdot \frac{\mu_0 I_1}{2\pi r}$$

$$\sin \alpha = \frac{r}{db}$$

$$\circ db = dz$$

$$\sin \alpha = \frac{r}{bz} \Rightarrow \sin \alpha \cdot bz = r$$

$$dr = \sin \alpha \cdot dz$$

$$dz = \frac{dr}{\sin \alpha}$$

$$F_{z1} = \int dF_{z1} = \int \frac{\mu_0 I_1}{2\pi r} I_2 dz = \int_{r_1}^{r_2} \frac{\mu_0 I_1}{2\pi r} I_2 \cdot \frac{dr}{\sin \alpha} = \frac{\mu_0 \cdot I_1 \cdot I_2}{2\pi \sin \alpha} \int_a^{a+b} \frac{dr}{r}$$

$$F_{z1} = \frac{\mu_0 \cdot I_1 \cdot I_2}{2\pi \sin \alpha} \cdot \ln \left(\frac{a+b}{a} \right)$$